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# Linear and nonlinear optical analysis on semiorganic *L*-proline cadmium chloride single crystal

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In the current investigation, *L*-proline cadmium chloride monohydrate (LPCC) single crystal is grown by a slow solvent evaporation technique to identify its credibility for nonlinear optical device applications. The constituent elements of LPCC crystal are determined by the energy dispersive spectroscopic (EDS) technique. The single crystal x-ray diffraction technique is used to determine the structural dimensions of LPCC crystal. The UV-visible studies are carried out within a wavelength range of 200 nm–1100 nm to determine the optical transmittance of LPCC crystal. The linear optical parameters of LPCC crystal are evaluated using the transmittance data to discuss its importance for distinct optical devices. The Nd:YAG laser assisted Kurtz–Perry test is carried out to determine the enhancement in second harmonic generation efficiency of LPCC crystal with reference to KDP crystal. The *Z*-scan technique is employed to assess the third order nonlinear optical (TONLO) properties of LPCC crystal at 632.8 nm. The *Z*-scan data are utilized to evaluate the TONLO refraction, absorption and susceptibility of LPCC crystal. The color oriented luminescence behavior of LPCC crystal is investigated within a spectral range of 350 nm–700 nm. The dependence of dielectric constant and dielectric loss on temperature and frequency is evaluated through the dielectric measurement studies.

Keywords: crystal growth, dielectric studies, nonlinear optical materials, optical studies

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### 1. Introduction

The challenge of designing efficient nonlinear optical (NLO) crystal has gained huge pace over few decades. Interestingly, the unexplored new class of semiorganic complex crystal is still an amateur but rapidly progressing field, which fascinates large audiences due to its bilateral contribution of properties associated with organic and inorganic counterparts. The coordination of organic and inorganic substrate reinforces acentric complex crystals, yielding a high nonlinear response.<sup>[1-4]</sup> In the current era of advanced functional materials it is a thought provoking fact that amino acid assisted organic-metal complex crystals stand out amongst the large regime of semiorganic NLO crystals are subject to an abundance of typical features such as high threshold to laser damage, large nonlinear response, low angular sensitivity, extended donor-acceptor network and high mechanical hardness, which fetches its desirability for practical utility in photonics, optoelectronics, optical modulation, frequency conversion, optical data storage and telecommunication devices.<sup>[5-8]</sup> As the presence of closed d<sup>10</sup> shell electrons in transition metals manifests a high optical response,<sup>[9]</sup> hence the growth of promising amino acid metal complex (AAMC) crystal has been an innovative idea which is currently on the verge and actively involves many research groups of the current era. Several potential AAMC crystals exhibiting an appreciable linear and nonlinear optical response are *L*-alanine sodium nitrate (LASN),<sup>[10]</sup> *L*-proline lithium bromide (LPLB),<sup>[11]</sup> L-alanine potassium chloride (LAPC),<sup>[12]</sup> L-valine cadmium chloride (LVCC),<sup>[13]</sup> bis L-proline cadmium iodide (BLPC),<sup>[14]</sup> L-alanine cadmium chloride,<sup>[15]</sup> Lproline magnesium chloride,<sup>[16]</sup> and L-proline cadmium chloride (LPCC).<sup>[17-20]</sup> The literature reveals that the LPCC crystal is an interesting NLO crystal which needs to be investigated thoroughly for identifying its credibility for distinct device applications. The optimization of growth parameters brings out significant changes in characteristic properties of crystal. The literature survey on the growth of LPCC crystal at different temperatures: 45 °C,<sup>[17]</sup> 27 °C,<sup>[19]</sup> 40 °C<sup>[20]</sup> reveals prominent changes in crystal growth, optical, dielectric, and mechanical properties. In the present study the growth tempera-

<sup>†</sup>Corresponding author. E-mail: loganees@gmail.com; loganees@yahoo.com © 2018 Chinese Physical Society and IOP Publishing Ltd ture of LPCC crystal is optimized to 36 °C while an additional novelty is that the complete third-order nonlinear optical properties of LPCC crystal are firstly investigated at 632.8 nm in this paper. Thus a unique approach is adopted to explore the features of LPCC crystal by employing crystal growth, single crystal x-ray diffraction, energy dispersive spectroscopy, UV-visible study, SHG efficiency test, Z-scan, dielectric and fluorescence characterization techniques.

#### 2. Experimental procedure

The AR grade raw material (L-proline and cadmium chloride) of purity 99% were purchased and recrystallized to eliminate the impurity at the primary stage of synthesis. The Lproline and cadmium chloride were later measured in stiochiometric ratio (1:1) and dissolved in double distilled water to allow the reaction for synthesis of L-proline cadmium chloride (LPCC) complex. The reaction mixture was allowed to agitate at a constant speed for six hours in order to facilitate homogeneous bonding. The reaction mixture was later filtered in a clean beaker through the membrane filter paper with the help of a vacuum pump. This beaker containing the filtered solution was covered with a porous sheet to avoid inclusion of dust particles and keenly placed in a constant temperature bath for slow evaporation at 36 °C. The seed crystals of LPCC were harvested within 10-12 days and the 21 mm  $\times$  10 mm  $\times$  0.7 mm bulk single crystal grown within a period of 20 days is shown in Fig. 1(a).



Fig. 1. (color online) (a) Single crystal of LPCC and (b) EDS spectrum of LPCC.

#### 3. Results and discussion

The LPCC crystal is subjected to the single crystal XRD technique at room temperature using the Enraf Nonius CAD4 crystal diffractometer. The experimentally determined structural parameters of LPCC crystal are discussed in Table 1. The analysis of the data reveals that LPCC crystallizes with orthorhombic crystal symmetry and belongs to the noncentrosymmetric space group  $P2_12_12_1$ . The structural parameters are in good agreement with those presented in Ref. [17].

Table 1. XRD data.		
Structural parameter	LPCC	LPCC <sup>[17]</sup>
Crystal system	Orthorhombic	Orthorhombic
Space group	P212121	P212121
a/Å	9.955	9.952
b/Å	13.479	13.484
c/Å	7.251	7.255
V/Å <sup>3</sup>	972.96	973.60

The elemental analysis of LPCC crystal is assessed by the energy dispersive spectroscopic (EDS) technique using the Hitachi S4700 instrument. The single crystal of LPCC is powdered and the energy spectrum is traced in an energy range of 0 keV–10 keV. The energy peaks corresponding to the elements are indexed in the spectrum shown in Fig. 1(b). The analysis of the EDS spectrum reveals the presence of carbon (47.43%), nitrogen (19.43%), oxygen (21.23%), chlorine (3.72%), and cadmium (8.19%). This evidences the formation of LPCC crystal complex.

Optical analysis gives the clues of associated electronic band structures and compositional nature of the material.<sup>[21,22]</sup> In the case of linear transmittance, the promotion of electrons to permitted energy states due to the absorption of incident photon and the inherent optical behavior of optically active functional units of the material are the principal parameters that determine the operative range of material in the optical spectrum.<sup>[23-25]</sup> In addition, constraints associated with single crystal such as anisotropy in molecular alignment along the crystal plane and defects (voids, vacancies, grain boundaries, striations, pits, solvent impurities, inclusions) tune the intensity of optical transmittance.<sup>[26-28]</sup> From the point of view of applications, the optical transmittances of NLO material in three prominent regions of spectrum far-UV (< 200 nm), visible (350 nm-700 nm) and NIR (ca. 1200 nm) need monitoring.<sup>[29]</sup> In the present case the LPCC crystal (2mm thickness) is scanned in a range of 200 nm-1100 nm using the spectrophotometer (Shimadzu-1601). The transmittance spectrum shown in Fig. 2(a) reveals that the LPCC crystal offers 84% transmittance which is uniform throughout the spectral range and the cutoff wavelength ( $\lambda_{cutoff}$ ) is found to be at 232 nm. The observed high transmittance in

LPCC crystal might have been governed by the low absorption tendency of amino acid (L-proline in the present case) and the minimum density of defects which sustain the impediment of free propagation of light in a crystal medium. The LPCC crystal with high transmittance within 240 nm-1100 nm claims its potential candidature for efficient transmission of the second and third harmonic signal of Nd:YAG laser operative at 1064 nm.<sup>[30]</sup> The transmittance of presently grown LPCC crystal is much superior to LPCC crystals grown at 45, 27, and 40 °C having transmittances of  $\sim 80\%$ ,  $\sim 60\%$ , and  $\sim 50\%$  respectively. The evaluated optical energy band gap  $(E_{\rm g} = 1243/\lambda_{\rm cutoff})^{[31]}$  of LPCC crystal turns out to be 5.35 eV. The thorough knowledge of linear optical parameters of the single crystal is essential to scrutinize the material for distinct optical applications. Hence for the variations of optical conductivity ( $\sigma_{op}$ , Fig. (2b)), extinction coefficient (K, Fig. 2(b)), refractive index (n, Fig. 2(c)), and reflectance (R, Fig. 2(b)) in the studied wavelength range are successfully evaluated using the formulae available in Ref. [32]. The analysis of graphs reveals that the optical conductivity and extinction coefficient increase with incident photon energy increasing. The LPCC crystal with high optical conductivity of the order of  $10^{13}$  and extinction coefficient of the order of  $10^{-5}$ possesses prerequisite qualities desirable for designing optical information processing and computing devices.<sup>[33]</sup> The resistance to the propagating light wave in a medium causes the velocity of light to change, which gives birth to the vital optical parameter refractive index. The refractive index of LPCC crystal is significantly lower in a range of 300 nm-1100 nm and materials with lower refractive index are in high demand for fabricating optical memory storage media for high resolution holographic systems<sup>[34]</sup> and calibrating the optical components (reflectors/filters/reflectors) used in solid state lasers.<sup>[35]</sup> The solar thermal device needs the coating material that can trap more and reflect less light incident on it and therefore the lower magnitude of reflectance throughout the visible region makes it a potential candidate as antireflecting coating material for solar thermal devices.<sup>[36]</sup>



Fig. 2. (color online) (a) UV-visible transmittance spectrum, (b) variation of optical conductivity and extinction coefficient, (c) variation of refractive index and reflectance, and (d) graphical representation of SHG response.

Crystals with high conversion efficiency are largely demanded by the photonic industry, hence the occurrence of the frequency doubling phenomenon in LPCC crystal has been experimentally determined using the Kurtz–Perry powder SHG efficiency test.<sup>[37]</sup> For sampling, all the components are aligned and the Nd:YAG laser (1064 nm, 8 ns, 10 Hz) is tuned to the Q-switched mode. The powder samples of LPCC and potassium dihydrogen phosphate (KDP) material are prepared and tightly packed in the microcapillary tube. Each sample is independently irradiated multiple times by the Gaussian filtered beam of laser and the output window is monitored. The emergence of green light is confirmed and the output signals are collimated through an array of photomultiplier tubes and further converted into an electrical signal to be displayed on an oscilloscope. The recorded voltages for KDP and LPCC crystal are shown in Fig. 2(d). The SHG efficiency of LPCC crystal is thus found to be 2.12 times higher than that of KDP. The high magnitude of SHG efficiency of LPCC crystal might be due to (i) the fact that the structural disordering/discrepancy in semiorganic molecular crystals leads to the formation of a sub-energy band below the energy gap, which increases the electron-phonon interaction and thus facilitates the favorable surrounding for photoinduced nonlinear optical effects to dwell efficiently,<sup>[38]</sup> and (ii) the enhanced charge transfer induced by the metal-organic coordination (donoracceptor bridge). The SHG efficiency of LPCC crystal is higher than LPLB (0.3 times that of KDP),<sup>[11]</sup> LPZC (0.5 times that of KDP), and BPCI (twice that of KDP)<sup>[14]</sup> crystals. Thus LPCC crystal with high SHG efficiency is a competent alternative for designing frequency convertors and SHG device elements.<sup>[39,40]</sup>

The focused beam irradiation of high intensity on the crystal surface dwells in an anharmonic manner in energy distribution due to which several third- order nonlinear optical effects evolve. The third order nonlinear optical (TONLO) effects are very sensitive and therefore accountability of the *Z*-scan technique is much higher for examining these TONLO properties of the given substrate. The *Z*-scan technique was proposed first by Bahae *et al.*<sup>[41]</sup> The configuration of *Z*-scan setup used for analysis is detailed in Table 2.

Table 2. The optical resolution of Z-scan setup.

Parameters and notations	Details
Laser wavelength $\lambda$ /nm	632.8
Lens focal length <i>f</i> /mm	30
Optical path distance Z/cm	85
Beam waist radius $\omega_a/mm$	3.3
Aperture radius $r_a$ /mm	2
Incident intensity at the focus $I_0/KW/m^2$	2.3375

The trace of TONLO refraction  $(n_2)$  in LPCC crystal is trailed using the closed aperture Z-scan analysis. To start the analysis, the He–Ne laser was optically aligned and the Gaussian filtered laser beam was converged via convex focusing lens to make localized irradiation on the crystal sample positioned at focus (Z = 0). The crystal is translated along the beam irradiated path, i.e., the Z direction which is equally divided into two halves with focus as the midpoint. The contributed change in transmitted light about the focus is monitored through the photo detector placed at the far field. The on-axis phase shift in TONLO refraction is governed by spatial distribution of energy along the crystal surface due to a localized thermal lensing effect<sup>[42]</sup> and it can be effectively tuned by modifying the repetition rate of the incident beam.<sup>[43,44]</sup> The closed aperture Z-scan transmittance curve shown in Fig. 3(a) depicts that the pre-focus peak followed by post-focus valley resembles the feature of negative nonlinear refraction which is the inherent characteristic of material possessing self-defocusing tendency.<sup>[45,46]</sup> The LPCC crystal with self-defocusing nature could be adopted for designing optical night vision sensor devices.<sup>[47]</sup> The peak-to-valley transmittance ( $\Delta T_{p-v}$ ) and on-axis phase shift ( $\Delta \phi$ ) are related by<sup>[41]</sup>

$$\Delta T_{p-\nu} = 0.406 (1-S)^{0.25} \left| \Delta \phi \right|, \tag{1}$$

where  $S = [1 - \exp(-2r_a^2/\omega_a^2)]$  is the aperture linear transmittance,  $r_a$  is the aperture radius, and  $\omega_a$  is the beam radius at the aperture. The value of  $n_2$  is evaluated using the following equation<sup>[41]</sup>

$$n_2 = \frac{\Delta \phi}{K I_0 L_{\rm eff}},\tag{2}$$

where  $K = 2\pi/\lambda$  ( $\lambda$  is the laser wavelength), *I* is the beam intensity at the focus,  $L_{\rm eff} = [1 - \exp(-\alpha L)]/\alpha$  is the effective thickness of the sample depending on linear absorption coefficient ( $\alpha$ ), and L is the thickness of the sample. The  $n_2$  of LPCC crystal is found to be  $6.3 \times 10^{-11}$  cm<sup>2</sup>/W. The TONLO absorption coefficient ( $\beta$ ) is examined and evaluated using the open aperture Z-scan analysis. The open aperture Z-scan curve of LPCC crystal is shown in Fig. 3(b). It reveals that as the sample approaches to the focus point the intensity of transmitted light becomes low, which is the concrete evidence for the reverse saturable absorption (RSA) effect.<sup>[48]</sup> In order to elaborate it, it must be realized that on a nanosecond time scale the existence of multiphoton absorption (MPA) and excited state absorption (ESA) phenomenon effectively reinforces the cross section of excited state absorption to capitalize the ground state absorption, which is the principle reason for the origin of RSA.<sup>[49,50]</sup> The TONLO absorption coefficient ( $\beta$ ) of grown crystal is evaluated from the following equation<sup>[41]</sup>

$$\beta = \frac{2\sqrt{2}\Delta T}{I_0 L_{\rm eff}},\tag{3}$$

where  $\Delta T$  is the one valley value at the open aperture Zscan curve. The  $\beta$  of LPCC crystal is found to be 7.29 × 10<sup>-5</sup> cm/W. The LPCC crystal with promising RSA property pronounces its liability for optical limiting devices.<sup>[51,52]</sup> The TONLO susceptibility ( $\chi^3$ ) is calculated using the following equations<sup>[41]</sup>

$$\operatorname{Re} \boldsymbol{\chi}^{(3)} (\operatorname{esu}) = 10^{-4} (\varepsilon_0 C^2 n_0^2 n_2) / \pi (\operatorname{cm}^2/\mathrm{W}), \qquad (4)$$

Im 
$$\chi^{(3)}$$
 (esu) =  $10^{-2} (\epsilon_0 C^2 n_0^2 \lambda \beta) / 4\pi^2$  (cm/W), (5)

$$\chi^{(3)} = \sqrt{(\text{Re}\,\chi^{(3)})^2 + (\text{Im}\,\chi^{(3)})^2},\tag{6}$$

where  $\varepsilon$  is the vacuum permittivity, *n* is the linear refractive index of the sample, and *c* is the velocity of light in a vacuum. The knowledge of polarizing ability of crystal material at high power laser intensities is emphasized to be the most essential factor and the  $\chi^3$  is the parameter which gives the magnitude of polarizability in the material. The  $\chi^3$  of LPCC crystal is found to be  $2.39 \times 10^{-4}$  esu. The delocalization rate of charge transfer over the pi bonded network is the vital factor giving rise to large  $\chi^3$ .<sup>[53]</sup>



**Fig. 3.** (color online) *Z*-scan transmittance curve with aperture (a) closed and (b) open, (c) PL emission spectrum.

Photoluminescence (PL) is the most vital nondestructive technique that sufficiently unfolds the electronic band structures, transition tendency of electrons in different energy states, electronic purity of the material, and the disorders associated with alloys.<sup>[54,55]</sup> The luminescence process involves the photo-excitation followed by photo-relaxation. The fluorophore (moiety of molecule responsible for its color) associated with the material is allowed to absorb the energy at the selected wavelength and re-emit the energy at a different but specific wavelength. The wavelength and intensity of emission relies on two factors: (i) the intrinsic nature of fluorophore and (ii) the surrounding environment of the fluorophore.<sup>[56]</sup> The surrounding environment implicates the presence of dopant/defects/impurities, which majorly affects the trajectory of an electron during photo-relaxation.<sup>[57]</sup> These aspects vitalize luminescence as a decisive feature which uncovers the exact property of the material. In the present analysis the luminescence study has been accomplished using the Hitachi FL-7000 spectrophotometer. Before the analysis, the instrument parameters were set to be excitation and emission slit width = 2.5 nm, scan speed = 240 nm/s, response time = 0.1 s, and delay = zero. The LPCC crystal material is photoexcited at an energy wavelength of 294 nm (4.22 eV) and the PL emission spectrum is recorded in a range of 350 nm-700 nm as shown in Fig. 3(c). It is observed that the LPCC crystal exhibits violet colored emission with high intensity and the peak maximum is centered at 421 nm. It is a known fact that the defects and impurities existing in a crystal avail more states for electrons and holes, which finally affects the transition energy, trajectory and lifetime of an electron.<sup>[58]</sup> The absence of multiple peaks, smooth emission profile and high intensity of emission confirms that the LPCC crystal does not possess optically active defects and intermediate states. In the fields of biomedical, biochemical, and chemical research, the materials with unique emission behaviors are utilized to detect the compounds, <sup>[59,60]</sup> and the LPCC crystal seems to meet the requirements for the discussed applications.

The dielectric measurement studies provide a great deal of information about the lattice dynamics and electric field distribution in material for exploiting the NLO applications.<sup>[61]</sup> The dielectric constant and dielectric loss of LPCC crystal are evaluated at different temperatures (30 °C-80 °C) and frequencies (100 Hz, 1 kHz, 10 kHz, 100 kHz, 1 MHz) using the HIOKI-3532 LCR cubemeter. The response of the polarization mechanism (electronic, ionic, dipolar, and space charge) is the principle factor responsible for the dielectric constant in crystal material.<sup>[62]</sup> The variation of the dielectric constant is shown in Fig. 4(a) and it is interesting to know that the dielectric constant of LPCC crystal increases with the increase in temperature but significantly decreases with the increase in frequency. It is due to the fact that all four polarizations are dominant at lower frequency and high temperature, which results in high magnitude of dielectric constant, on the contrary, the decrease in magnitude of dielectric constant is less at higher frequencies due to the diminished effect of foresaid polarizations.<sup>[63]</sup> The space charge polarization plays an active role as the temperature of the material increases but as the frequency of applied field exceeds 1 kHz even the space charge



Fig. 4. (color online) Response of (a) dielectric constant and (b) dielectric loss.

polarization becomes deficient.<sup>[64]</sup> It is an intruding observation that the LPCC crystal is reported to have a high dielectric constant in a given frequency range,<sup>[19,65]</sup> however much lower, and the concurring results are obtained in the presently grown LPCC crystal, which might have stable growth at constrained temperature, high electronic purity and less hindrance of space charge polarization. It is highlighted that the dielectric constant of LPCC crystal has been tuned to much lower value as frequency increased to 1 MHz. The materials exhibiting low dielectric constant consume less power and offer decreased RC delay, which is vital for microelectronics devices<sup>[66]</sup> and also enhances the SHG coefficient of the material.<sup>[67]</sup> Thus the LPCC crystal with lower dielectric constant is considered to be a better alternative for fabricating photonics, electro-optic modulators, field detectors, THz wave generators, and microelectronic devices.<sup>[68,69]</sup> The dissipation of electromagnetic energy in crystal medium is termed dielectric loss. The dielectric loss in single crystal is majorly contributed by macro/micro cracks, random crystal orientation, porosity, impurities, grain boundaries, etc.<sup>[70,71]</sup> The response of dielectric loss shown in Fig. 4(b) reveals that the dielectric loss follows the same trend as that of the dielectric constant. The lower magnitude of dielectric loss at higher frequencies suggests that the LPCC crystal possesses a very small quantity of electrically active defects.<sup>[72]</sup> It is noteworthy that the dielectric constant and dielectric loss show a lot less variation at higher frequencies even when the temperature of the crystal sample is raised. From an application point of view the LPCC crystal could be a potential candidate for distinct technological devices.

#### 4. Conclusions

LPCC bulk single crystal is successfully grown by the slow solvent evaporation technique. The EDS analysis confirmed the presence of constituent elements of LPCC crystal. In single crystal XRD analysis the LPCC crystal is found to have an orthorhombic crystal structure with space group P2<sub>1</sub>2<sub>1</sub>2<sub>1</sub>. Linear optical studies employed within 200 nm-1100 nm reveal that the LPCC crystal has high optical transparency (84%), low cutoff wavelength (232 nm), wide optical band gap (5.35 eV), low refractive index, reflectance, extinction coefficient and increasing optical conductivity. The promising second and TONLO property of LPCC crystal is confirmed by Kurtz-Perry test and Z-scan analysis. The SHG efficiency of LPCC crystal is found to be 2.12 times higher than that of KDP crystal. The LPCC crystal possesses the negative nonlinear refraction and reverse saturable absorption under beam irradiation of wavelength 632.8 nm. The values of  $n_2$ ,  $\beta$ , and  $\chi^3$  of LPCC crystal are  $6.3 \times 10^{-11}$  cm<sup>2</sup>/W,  $7.29 \times 10^{-5}$  cm/W,  $2.39 \times 10^{-4}$  esu, respectively. The luminescence behavior investigated with 350 nm-700 nm confirms the prominence of violet colored emission in LPCC crystal with a maximum at 421 nm. The dielectric studies reveal that the dielectric constant and dielectric loss of LPCC crystal greatly decreases with the increase in frequency throughout the applied temperature range of 30 °C-80 °C. The LPCC crystal with high optical transparency, large SHG efficiency, promising TONLO properties, single colored emission, and low dielectrics advocates the efficient utility of LPCC crystal for applications in photonics, NLO, night vision sensors, biomedical, optical limiting, SHG, and microelectronics devices.

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#### References

[1] Wang X Q, Xu D, Lu M K, Yuan D R and Xu S X 2001 Mater. Res. Bull. 36 879

- [2] He Y H, Lan Y Z, Zhan C H, Feng Y L and Su H 2009 Inorg. Chem. Acta 362 1952
- [3] Verbiest T, Houbrechts S, Kauranen M, Clays K and Persoons A 1997 J. Mater. Chem. 7 2175
- [4] Lu G W, Xia H R, Wang X Q, Xu D, Chen Y and Zhou Y Q 2001 *Mater. Sci. Eng. B* 87 117
- [5] Sugandhi K, Verma S, Jose M, Joseph V and Das S J 2013 Opt. Laser Technol. 54 347
- [6] Peter M E and Ramasamy P 2010 J. Cryst. Growth 312 1952
- [7] Hanumantharao R and Kalainathan S 2012 Mater. Lett. 74 74
- [8] Liu X J, Xu D, Wei X Q, Ren M J and Zhang G H 2010 Mater. Sci. Eng. B 166 203
- [9] Bright K C and Freeda T H 2010 Physica B 405 3857
- [10] Sethuraman K, Babu R R, Gopalakrishnan R and Ramasamy P 2008 Cryst. Growth Des. 8 1863
- [11] Shkir M, Alfaify S, Ajmal Khan M, Dieguez E and Perles J 2014 J. Cryst. Growth 391 104
- [12] Arockia Avila S, Selvakumar S, Francis M and Rajesh A L 2017 J. Mater. Sci. Mater. Electron. 28 1051
- [13] Maadeswaran P and Chandrasekaran J 2011 Optik 122 1128
- [14] Boopathi K, Jagan R and Ramasamy P 2016 Appl. Phys. A 122 689
- [15] Raj C J and Das S J 2008 Cryst. Growth Des. 8 2729
- [16] Anbuselvi D, Theras J E M, Jayaraman D and Joseph V 2013 Physica B 423 38
- [17] Kandasamy A, Siddeswaran R, Murugakoothan P, Kumar P S and Mohan R 2007 Cryst. Growth Des. 7 183
- [18] Prakash J T J and Kumararaman S 2008 Mater. Lett. 62 4097
- [19] Shakir M, Kushwaha S K, Maurya K K, Bhatt R C, Rashmi, Wahab M A and Bhagavannarayana G 2010 Mater. Chem. Phys. 120 566
- [20] Singh P, Hasmuddin M, Shakir M, Vijayan N, Abdullah M M, Ganesh V and Wahab M A 2013 Mater. Chem. Phys. 142 154
- [21] Sun Z H, Xu D, Wang X Q, Zhang G H, Yu G, Zhu L Y and Fan H L 2009 Mater. Res. Bull. 44 925
- [22] Sivakumar N, Kanagathara N, Varghese B, Bhagavannarayana G, Gunasekaran S and Anbalagan G 2014 Spectrochem. Acta A 118 603
- [23] Shaikh R N, Anis M, Shirsat M D and Hussaini S S 2014 J. Optoelectron. Adv. Mater. 16 1147
- [24] Kakani S L and Kakani A 2004 Material Science (New Delhi: New Age International) p. 417
- [25] Anis M, Ramteke S P, Shirsat M D, Muley G G and Baig M I 2017 Opt. Mater. 72 590
- [26] Weber M J 2003 Handbook of Optical Materials (New York: CRC Press)
- [27] Bass M 1995 Handbook of Optics, Volume 2 (USA: Mc-Graw Hill) p. 33.12
- [28] Ramteke S P, Baig M I, Shkir M, Kalainathan S, Shirsat M D, Muley G G and Anis M 2018 Opt. Laser Technol. 104 83
- [29] Ganapayya B, Jayarama A, Sankolli R, Hathwar V R and Dharmaprakash S M 2012 J. Mol. Struc. 1007 175
- [30] Anis M, Pandian M S, Baig M I, Ramasamy P and Muley G G 2017 Mater. Res. Innov.
- [31] Sangeetha K, Prasad L G and Mathammal R 2016 Physica B 501 5
- [32] Rajagopalan N R and Krishnamoorthy P 2017 J. Inorg. Organomet. Polym. Mater. 27 296
- [33] Anis M and Muley G G 2016 Phys. Scr. 91 85801
- [34] Chen C T and Liu G Z 1986 Ann. Rev. Mater. Sci. 16 203
- [35] Schubert E F, Kim J K and Xi J Q 2007 Phys. Stat. Sol. B 244 3002
- [36] Raja M V A, Anand D P and Madhavan J 2013 Sci. ActaXaveriana 4 41
- [37] Kurtz S K and Perry T T 1968 J. Appl. Phys. 39 3698
- [38] Wojciechowski A, Alzayed N, Kityk I, Berdowski J, and Tylczyńki Z 2010 Opt. Appl. XL 1007

- [39] Pahurkar V G, Anis M, Baig M I, Ramteke S P and Muley G G 2017 Optik 142 421
- [40] Li Y, Wang M, Zhu T, Meng X, Zhong C, Chen X and Qin J 2012 Dalton Trans. 41 763
- [41] Bahae M S, Said A A, Wei T H, Hagan D J and Stryland E W V 1990 IEEE J. Quantum Electron. 26 760
- [42] Baig M I, Anis M, Kalainathan S, Babu B and Muley G G 2017 Mater. Technol. Adv. Perform. Mater. 32 560
- [43] Gu B, Wang H T and Ji W 2009 Opt. Lett. 34 2769
- [44] Ramteke S P, Anis M, Pandian M S, Kalainathan S, Baig M I, Ramasamy P and Muley G G 2018 Opt. Laser Technol. 99 197
- [45] Sun X B, Wang X Q, Ren Q, Zhang G H, Yang H L and Feng L 2006 Mater. Res. Bull. 41 177
- [46] Rasal Y B, Anis M, Shirsat M D and Hussaini S S 2017 Mater. Res. Innov.
- [47] Muthuraja A and Kalainathan S 2017 Mater. Technol. Adv. Perform. Mater. 32 335
- [48] Azhar S M, Anis M, Hussaini S S, Kalainathan S, Shirsat M D and Rabbani G 2016 Mater. Sci. Poland 34 800
- [49] Zhang Y D, Zhao Z Y, Yao C B, Yang L, Li J and Yuan P 2014 Opt. Laser Technol. 58 207
- [50] Zidana M D, Al-Ktaifani M M and Allahham A 2017 Opt. Laser Technol. 90 174
- [51] Azhar S M, Hussaini S S, Shirsat M D, Rabbani G, Shkir M, AlFaify S, Ghramh H A, Baig M I and Anis M 2017 Mater. Res. Innov.
- [52] Kumar R S S, Rao S V, Giribabu L and Rao D N 2007 Chem. Phys. Lett. 447 274
- [53] Shaikh R N, Anis M, Shirsat M D and Hussaini S S 2018 Optik 154 435
- [54] Anis M, Hussaini S S, Shirsat M D, Shaikh R N and Muley G G 2016 Mater. Res. Express 3 106204
- [55] Anis M, Baig M I, Pandian M S, Ramasamy P, AlFaify S, Ganesh V, Muley G G and Ghramh H A 2018 Cryst. Res. Technol.
- [56] Sauer M, Hofkens J and Enderlein J 2011 Handbook of Fluorescence Spectroscopy and Imaging (Weinheim: Wiley-VCH)
- [57] Chandran S, Paulraj R and Ramasamy P 2015 Mater. Res. Bull. 68 210
- [58] Timothy H G 2000 Photoluminescence in analysis of surfaces and interfaces, in Encyclopedia of Analytical Chemistry (Chichester: John Wiley & Sons) p. 9209
- [59] Kajamuhideen M S, Senthuraman K, Ramamurthi K and Ramasamy P 2017 Opt. Laser Technol. 91 159
- [60] Shaikh R N, Anis M, Shirsat M D and Hussaini S S 2015 Mater. Technol. Adv. Perform. Mater. 19 187
- [61] Subhashini R, Sathya D, Sivashankar V, Mageshwari P S L and Arjunan S 2016 Opt. Mater. 62 357
- [62] Riscob B, Shakir M, Sundar J K, Natarajan S, Wahab M A and Bhagavannarayana G 2011 Spectrochem. Acta A 78 543
- [63] Anis M, Muley G G, Baig M I, Hussaini S S and Shirsat M D 2017 Mater. Res. Innov. 21 439
- [64] Singh P, Hasmuddin M, Abdullah M M, Shkir M and Wahab M A 2013 Mater. Res. Bull. 48 3926
- [65] Kandasamy A, Mohan R, Caroline M L and Vasudevan S 2008 Cryst. Res. Technol. 43 186
- [66] Hatton B D, Landskron K, Hunks W J, Bennett M R, Shukaris D, Perovic D D and Ozin G A 2006 Mater. Today 9 22
- [67] Miller R C 1964 Appl. Phys. Lett. 5 17
- [68] Jazbinsek M, Mutter L and Gunter P 2008 IEEE J. Sel. Top. Quantum Electron. 14 1298
- [69] Anis M, Muley G G, Pahurkar V G, Baig M I and Dagdale S R 2018 Mater. Res. Innov. 22 99
- [70] Jiang M H and Fang Q 1999 Adv Mater. 11 1147
- [71] Baig M I, Anis M and Muley G G 2017 Opt. Mater. 72 1
- [72] RamtekeS P, Anis M, Baig M I and Muley G G 2018 Optik 154 275